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# Existence, Uniqueness and Structure of Second Order Absolute Minimisers

by

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#### EXISTENCE, UNIQUENESS AND STRUCTURE OF SECOND ORDER ABSOLUTE MINIMISERS

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ABSTRACT. Let  $\Omega \subseteq \mathbb{R}^n$  be a bounded open  $C^{1,1}$  set. In this paper we prove the existence of a unique second order absolute minimiser  $u_\infty$  of the functional

$$E_{\infty}(u, \mathcal{O}) := \|F(\cdot, \Delta u)\|_{L^{\infty}(\mathcal{O})}, \quad \mathcal{O} \subseteq \Omega \text{ measurable},$$

with prescribed boundary conditions for u and Du on  $\partial\Omega$  and under natural assumptions on F. We also show that  $u_{\infty}$  is partially smooth and there exists a harmonic function  $f_{\infty} \in L^{1}(\Omega)$  such that

$$F(x, \Delta u_{\infty}(x)) = e_{\infty} \operatorname{sgn}(f_{\infty}(x))$$

for all  $x \in \{f_{\infty} \neq 0\}$ , where  $e_{\infty}$  is the infimum of the global energy.

#### 1. INTRODUCTION

For  $n \in \mathbb{N}$ , let  $\Omega \subseteq \mathbb{R}^n$  be a bounded open set and let also  $F : \Omega \times \mathbb{R} \longrightarrow \mathbb{R}$  be a real function that is  $\mathcal{L}(\Omega) \otimes \mathcal{B}(\mathbb{R})$ -measurable, namely, measurable with respect to the product  $\sigma$ -algebra of the Lebesgue subsets of  $\Omega$  with the Borel subsets of  $\mathbb{R}$ . In this paper we consider variational problems for second order supremal functionals of the form

(1.1) 
$$\operatorname{E}_{\infty}(u, \mathcal{O}) := \| \operatorname{F}(\cdot, \Delta u) \|_{L^{\infty}(\mathcal{O})}, \quad \mathcal{O} \subseteq \Omega \text{ measurable},$$

where the admissible functions u range over the (Fréchet) Sobolev space

(1.2) 
$$\mathcal{W}^{2,\infty}(\Omega) := \bigcap_{1$$

The following is a natural notion of minimiser for variational problems of this type.

**Definition 1** (Second order absolute minimisers). A function  $u \in W^{2,\infty}(\Omega)$  is called a second order absolute minimiser of (1.1) on  $\Omega$  if

$$\mathbf{E}_{\infty}(u, \mathcal{O}) \leq \mathbf{E}_{\infty}(u + \phi, \mathcal{O})$$

for all open sets  $\mathcal{O} \subseteq \Omega$  and all  $\phi \in \mathcal{W}^{2,\infty}_0(\mathcal{O})$ .

Here we have used the obvious notation  $\mathcal{W}_0^{2,\infty}(\mathcal{O}) := \mathcal{W}^{2,\infty}(\mathcal{O}) \cap \mathcal{W}_0^{2,2}(\mathcal{O})$ . Given  $u_0 \in \mathcal{W}^{2,\infty}(\Omega)$ , we will also write  $\mathcal{W}_{u_0}^{2,\infty}(\Omega) = u_0 + \mathcal{W}_0^{2,\infty}(\Omega)$ .

The concept of an absolute minimiser would be unnecessary for variational problems given in terms of integrals rather than the essential supremum, where it would suffice to consider global minimisers on  $\Omega$  for fixed boundary data. But here, since the functional (1.1) is not a measure (not additive) in the domain argument, a global minimiser of  $E_{\infty}(\cdot, \Omega)$  will not necessarily be a second order absolute minimiser.

Key words and phrases.  $\infty$ -Laplacian;  $\infty$ -Bilaplacian; Second Order absolute minimisers; Calculus of Variations in  $L^{\infty}$ .

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**Example.** For  $\Omega = (-1,0) \cup (0,1)$ , consider the functional  $E_{\infty}(u, \mathcal{O}) = ||u''||_{L^{\infty}(\mathcal{O})}$ and let  $Q(x) := x\chi_{(-1,0)}(x) - x(x-1)\chi_{(0,1)}(x)$ ,  $x \in \Omega$ . Then Q is a global minimiser of  $E_{\infty}(\cdot, \Omega)$  in  $\mathcal{W}_Q^{2,\infty}(\Omega)$  with  $E_{\infty}(Q,\Omega) = 2$  and is also a second order absolute minimiser. However, for any function  $\zeta \in C_c^{\infty}(-1,0)$  with  $0 < ||\zeta''||_{L^{\infty}(-1,0)} < 1$ , the perturbation  $Q + \zeta$  still satisfies  $E_{\infty}(Q + \zeta, \Omega) = 2$  and lies in  $\mathcal{W}_Q^{2,\infty}(\Omega)$ , but does not minimise  $E_{\infty}(\cdot, (-1,0))$  over  $W_Q^{2,\infty}((-1,0))$  because the only minimiser on (-1,0) with boundary data Q is the identity.

On the other hand, if  $u \in W^{2,\infty}(\Omega)$  minimises  $E_{\infty}(\cdot, \Omega)$  uniquely in  $W^{2,\infty}(\Omega)$  with respect to its own boundary conditions, then we will show that u is actually the unique second order absolute minimiser. This is the situation that we will find in the main results of this paper. We first give a condition that guarantees that u is the unique minimiser for its boundary values on any subdomain. While it is not obvious that this condition can be met, we will subsequently prove that it is satisfied by exactly one function under given boundary conditions and very mild additional assumptions.

In the following, we will use the symbolisation "sgn" for the sign function, with the convention that sgn(0) = 0. We will assume also that

(1.3) 
$$\Omega$$
 is a bounded connected open subset of  $\mathbb{R}^n$ ,  $n \ge 1$ ,

and

(1.4) 
$$\begin{cases} F: \Omega \times \mathbb{R} \longrightarrow \mathbb{R} \text{ is } \mathcal{L}(\Omega) \otimes \mathcal{B}(\mathbb{R}) \text{-measurable and for a.e.} \\ x \in \Omega, \xi \mapsto F(x,\xi) \text{ is strictly increasing with } F(x,0) = 0. \end{cases}$$

Our first main result therefore is:

**Theorem 2** (Criterion for unique minimisers). Suppose (1.3)-(1.4) hold and consider (1.1) and a function  $u_* \in W^{2,\infty}(\Omega)$ . If there exist a number  $e_* \geq 0$  and a function  $f_* \in L^1(\Omega)$  satisfying

(1.5) 
$$\Delta f_* = 0, \quad on \ \Omega$$

such that

(1.6) 
$$\mathbf{F}(\cdot, \Delta u_*) = e_* \operatorname{sgn}(f_*), \quad a.e. \text{ on } \Omega,$$

then

$$\mathrm{E}_{\infty}(u_*,\mathcal{O}) < \mathrm{E}_{\infty}(u_*+\phi,\mathcal{O}),$$

for any open  $\mathcal{O} \subseteq \Omega$  and any  $\phi \in \mathcal{W}^{2,\infty}_0(\mathcal{O}) \setminus \{0\}$ .

Namely, for any open subset  $\mathcal{O} \subseteq \Omega$ , the function  $u_*$  is the unique global minimiser of  $E_{\infty}(\cdot, \mathcal{O})$  in  $\mathcal{W}^{2,\infty}_{u_*}(\mathcal{O})$ .

Consequently, any such function  $u_*$  as in Theorem 2 is the unique second order absolute minimiser as well as the unique minimiser in  $\mathcal{W}^{2,\infty}(\Omega)$  with respect to its own boundary conditions on  $\partial\Omega$ .

Clearly, (1.6) implies that  $e_* = \mathcal{E}_{\infty}(u_*, \Omega)$  unless  $f_* \equiv 0$ . If we fix boundary data by choosing  $u_0 \in \mathcal{W}^{2,\infty}(\Omega)$  and insist that  $u_* \in \mathcal{W}^{2,\infty}_{u_0}(\Omega)$ , then there is at most one minimiser  $u_* \in \mathcal{W}^{2,\infty}_{u_0}(\Omega)$  that satisfies the condition and  $e_*$  is uniquely determined as the infimum of  $\mathcal{E}_{\infty}(\cdot, \Omega)$  over the space. On the other hand, the function  $f_*$  is not determined uniquely by the boundary conditions. Note further that since  $\mathcal{F}(x, \cdot)$  is by assumption strictly increasing, (1.6) is equivalent to the next representation formula for  $u_*$ :

$$\Delta u_*(x) = \mathbf{F}(x, \cdot)^{-1} \Big( e_* \operatorname{sgn} \big( f_*(x) \big) \Big), \quad \text{a.e. } x \in \Omega.$$

Moreover, by using standard argument involving Green functions (see e.g. [GT, Ch. 2]), we could represent  $u_*$  in terms of F,  $f_*, e_*, u_*|_{\partial\Omega}, D_{\nu}u_*|_{\partial\Omega}$ .

Theorem 2 gives a connection between the variational problem and a PDE system of second order equations with a parameter consisting of (1.5) and (1.6). We will see later that under certain assumptions on  $\Omega$ , F and the boundary data, the system has in fact a solution  $(u_*, f_*, e_*)$  with  $f_* \neq 0$  if  $e_* > 0$ . It then follows that  $(u_*, e_*)$ is unique and the system is equivalent to unique global minimality under prescribed boundary data. We may think of (1.5) and (1.6) as a PDE formulation of the  $L^{\infty}$ variational problem. There does exist, however, a more conventional analogue of the "Euler-Lagrange equation" for (1.1). This is the fully nonlinear PDE of third order

(1.7) 
$$\mathbf{F}(\cdot, \Delta u) \mathbf{F}_{\xi}(\cdot, \Delta u) \left| \mathbf{D} \left( |\mathbf{F}(\cdot, \Delta u)|^2 \right) \right|^2 = 0, \quad \text{on } \Omega.$$

A particular model case of (1.1)-(1.7) is what we call the " $\infty$ -Bilaplacian" and arises from the choice  $F(x,\xi) = \xi$ . Then, equation (1.6) becomes  $\Delta u_* = e_* \operatorname{sgn}(f_*)$  and (1.7) becomes  $\Delta u |D(|\Delta u|^2)|^2 = 0$ . Due to the particular structure of the functional (1.1), however, in this case (1.7) becomes redundant since all the structural information of second order absolute minimisers can be obtained directly from the  $L^{\infty}$  variational problem.

For our existence result, we will assume that  $F: \Omega \times \mathbb{R} \longrightarrow \mathbb{R}$  satisfies

(1.8) 
$$\begin{cases} F \in C^2(\Omega \times \mathbb{R}), \\ F(x,0) = 0, \quad x \in \Omega, \\ \exists c > 0 : \begin{cases} c \leq F_{\xi}(x,\xi) \leq \frac{1}{c}, \quad (x,\xi) \in \Omega \times \mathbb{R}, \\ F(x,\xi) F_{\xi\xi}(x,\xi) \geq -\frac{1}{c}, \quad (x,\xi) \in \Omega \times \mathbb{R}, \end{cases}$$

where subscripts of F denote partial derivatives. The conditions of (1.8) imply that for any fixed  $x \in \Omega$  the partial function  $|F(x, \cdot)|$  is level-convex on  $\mathbb{R}$  (i.e. has convex sublevel sets) but in general they do not imply convexity.

In addition to existence, we will prove that the absolute minimisers of the  $L^{\infty}$  problem can be approximated by solutions of corresponding minimisers of  $L^p$  problems as  $p \to \infty$ . For  $p \in (1, \infty)$ , we therefore define

$$\mathbf{E}_{p}(u) := \left( \oint_{\Omega} \left| \mathbf{F}(\cdot, \Delta u) \right|^{p} \right)^{1/p}, \quad u \in W^{2, p}(\Omega)$$

where the slashed integral sign denotes the average over  $\Omega$ .

**Theorem 3** (Existence, structure and approximation). Let  $\Omega$  satisfy (1.3) and also have  $C^{1,1}$  boundary. Let also F satisfy (1.8) and fix a function  $u_0$  in  $W^{2,\infty}(\Omega)$  with  $\Delta u_0 \in C(\overline{\Omega})$ .

(I) There exist a global minimiser  $u_{\infty} \in \mathcal{W}^{2,\infty}_{u_0}(\Omega)$  of  $\mathcal{E}_{\infty}(\cdot, \Omega)$  and a harmonic function  $f_{\infty} \in L^1(\Omega)$  such that

$$F(\cdot, \Delta u_{\infty}) = e_{\infty} \operatorname{sgn}(f_{\infty}), \quad a.e. \text{ on } \Omega,$$

where  $e_{\infty} = E_{\infty}(u_{\infty}, \Omega)$ . Further,  $f_{\infty} \not\equiv 0$  if  $e_{\infty} > 0$ . (II) Let

$$\Gamma_{\infty} := f_{\infty}^{-1}(\{0\}).$$

If  $e_{\infty} > 0$ , then  $u_{\infty}$  belongs to  $C^{3,\alpha}(\Omega \setminus \Gamma_{\infty})$  for any  $\alpha \in (0,1)$  and  $\Gamma_{\infty}$  is a Lebesgue nullset. If  $e_{\infty} = 0$ , then  $u_{\infty}$  is a harmonic function.

(III) For  $p \in \mathbb{N}$ , let

$$e_p := \inf \{ E_p(u) : u \in W^{2,p}_{u_0}(\Omega) \}.$$

Then, for any p large enough there exists a global minimiser  $u_p \in W^{2,p}_{u_0}(\Omega)$ of  $E_p$  satisfying  $E_p(u_p) = e_p$ . Moreover,  $e_p \longrightarrow e_\infty$  as  $p \to \infty$  and there exists a subsequence  $(p_\ell)^{\infty}_{\ell=1}$  such that  $u_{p_\ell} \longrightarrow u_\infty$  in the weak topology of  $\mathcal{W}^{2,\infty}(\Omega)$  as  $\ell \to \infty$ . In addition,

$$\begin{array}{ll} \left\{ \begin{array}{ll} u_{p_{\ell}} \longrightarrow u_{\infty}, & \text{ in } C^{1}(\overline{\Omega}), \\ \mathrm{D}^{2} u_{p_{\ell}} \longrightarrow \mathrm{D}^{2} u_{\infty}, & \text{ in } L^{q}(\Omega, \mathbb{R}^{n \times n}) \text{ for all } q \in (1, \infty), \\ \Delta u_{p_{\ell}} \longrightarrow \Delta u_{\infty}, & \text{ a.e. on } \Omega \text{ and in } L^{q}(\Omega) \text{ for all } q \in (1, \infty), \end{array} \right.$$

as  $\ell \to \infty$ . Furthermore,  $\Delta u_{p_{\ell}} \longrightarrow \Delta u_{\infty}$  locally uniformly on  $\Omega \setminus \Gamma_{\infty}$  if  $e_{\infty} > 0$  and locally uniformly on  $\Omega$  if  $e_{\infty} = 0$ .

$$(IV)$$
 Let

(1.9) 
$$f_p := \frac{1}{e_p^{p-1}} \left| \mathbf{F}(\cdot, \Delta u_p) \right|^{p-2} \mathbf{F}(\cdot, \Delta u_p) \mathbf{F}_{\xi}(\cdot, \Delta u_p)$$

if  $e_p \neq 0$  and  $f_p \equiv 0$  if  $e_p = 0$ . Then, the harmonic function  $f_{\infty}$  in (I) may be chosen such that  $f_{p_{\ell}} \longrightarrow f_{\infty}$  as  $\ell \to \infty$  in the strong local topology of  $C^{\infty}(\Omega)$ .

By invoking Theorem 2, an immediate consequence is that the modes of convergence in Theorem 3(III) as  $p \to \infty$  are actually full and not just subsequential. Also, known results on the regularity of nodal sets of solutions to elliptic equations [HS] imply that  $\Gamma_{\infty}$  is countably rectifiable, being equal to the union of countably many smooth submanifolds of  $\Omega$  and a set of vanishing (n-2)-dimensional Hausdorff measure. This, however, uses only the fact that  $f_{\infty}$  is harmonic and it seems plausible that the full statement in (I) could give further information.

The optimal regularity of  $\Gamma_{\infty}$  is an open question which we do not attempt to answer here. Certainly, full regularity of the set  $\Gamma_{\infty}$  cannot be expected as there are limitations: in [KP2] it was noted that  $\Gamma_{\infty}$  in general may not be a smooth submanifold, as for certain data  $u_0$  the intersection of transversal lines in  $\Omega$  was observed in numerical experiments. In most cases the set  $\Gamma_{\infty}$  is necessarily nonempty and divides  $\Omega$  into two distinct parts, whilst the equations  $F(\cdot, \Delta u_{\infty}) = \pm e_{\infty}$ do not permit any solutions in  $\mathcal{W}^{2,\infty}_{u_0}(\Omega)$  for most boundary data, even if n = 1. Therefore, the Laplacian of  $u_{\infty}$  will have a jump on  $\Gamma_{\infty}$  and, in terms of  $u_{\infty}$ , no more regularity than  $\mathcal{W}^{2,\infty}(\Omega)$  can be expected (see the numerical and explicit solutions in [KP2]).

Let us also note further that (I)–(IV) above have been obtained in [KP2] for n = 1 and in some other special cases (although were not stated in this explicit fashion), whilst the qualitative behaviour emerging here was observed numerically for n = 2 and  $F(x, \xi) = \xi$ .

When combined, Theorem 2 and Theorem 3 imply in particular the following.

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**Corollary 4.** Under the hypotheses of Theorem 3, there exists a unique global minimiser  $u_{\infty}$  of  $E_{\infty}(\cdot, \Omega)$  in  $\mathcal{W}^{2,\infty}_{u_0}(\Omega)$ , which is a second order absolute minimiser and a strong solution to the Dirichlet problem for (1.7):

$$\begin{cases} F(\cdot, \Delta u)F_{\xi}(\cdot, \Delta u)|D(|F(\cdot, \Delta u)|^{2})|^{2} = 0, & \text{in } \Omega,\\ u = u_{0}, & \text{on } \partial\Omega,\\ Du = Du_{0}, & \text{on } \partial\Omega, \end{cases}$$

More precisely,  $u_{\infty}$  is thrice differentiable a.e. on  $\Omega$  and satisfies the PDE in the pointwise sense.

The study of supremal functionals and of their associated equations is known as Calculus of Variations in the space  $L^{\infty}$ . Second order variational problems in  $L^{\infty}$  have only relatively recently been studied and are still poorly understood. It is remarkable that for our specific problem, we obtain not just unique absolute minimisers, but also a fair amount of detailed information about their structures, with relatively simple means. On the other hand, our methods take advantage of the special structure of the problem and are unlikely to work in general, although they allow the following modest generalisation: all of the preceding results hold for the seemingly more general case where the Laplacian is replaced by the projection  $\mathbf{A}: D^2 u = \sum_{i,j} \mathbf{A}_{ij} D_{ij}^2 u$  on a fixed positive symmetric matrix  $\mathbf{A} \in \mathbb{R}^{n \times n}_+$ . This gives rise to the following functional:

$$\mathbf{E}_{\infty}(u, \mathcal{O}) = \left\| \mathbf{F}(\cdot, \mathbf{A} : \mathbf{D}^{2}u) \right\|_{L^{\infty}(\mathcal{O})}, \quad \mathcal{O} \subseteq \Omega \text{ measurable.}$$

However, this case can easily be reduced to the case we study herein via the change of variables  $x \mapsto \Lambda O^{\top} x$  for a diagonal  $n \times n$  matrix  $\Lambda$  and an orthogonal matrix  $O \in O(n)$  arising from the spectral representation  $\mathbf{A} = O\Lambda^2 O^{\top}$ .

Some of the techniques that underpin Theorem 3 have been successfully deployed to problems somewhat different to (1.1) (with dependence on u in addition to  $\Delta u$ ) [MS, S1], which suggests that further generalisation might be possible. In order to keep the presentation simple, however, we do not explore this possibility any further in this work.

We conclude this introduction by placing the  $L^{\infty}$  problem we study herein into the wider context of Calculus of Variations. Variational problems for first order functionals of the form

(1.10) 
$$(u, \mathcal{O}) \longmapsto \operatorname{ess\,sup}_{x \in \mathcal{O}} \operatorname{H}(x, u(x), \operatorname{D}u(x)), \quad u \in W^{1,\infty}(\Omega), \ \mathcal{O} \subseteq \Omega,$$

together with the associated PDEs, first arose in the work of Aronsson in the 1960s [A1]–[A3]. The first order case is very well developed and the relevant bibliography is very extensive. For a pedagogical introduction to the theme which is accessible to non-experts, we refer to the monograph [K8] (see also [C]). The vectorial case of (1.10) for maps  $u : \mathbb{R}^n \supseteq \Omega \longrightarrow \mathbb{R}^N$  is a much more modern and rapidly developing topic which first arose in recent work of the first author in the early 2010s (see [K1]–[K7], [K9]–[K13] as well as the joint works with Abugirda, Ayanbayev, Croce, Kristensen, Manfredi, Pisante and Pryer [AK, KM, CKP, KP, KK, AyK]).

In a very recent paper, the first author, jointly with Pryer (see [KP2] and also [KP3]), initiated the study of higher order variational problems and of their associated PDEs. As a first step they considered functionals of the form

$$(u, \mathcal{O}) \longmapsto \operatorname{ess\,sup}_{x \in \mathcal{O}} \mathrm{H}(\mathrm{D}^2 u(x)), \quad u \in W^{2,\infty}(\Omega), \ \mathcal{O} \subseteq \Omega,$$

with dependence on pure second derivatives only. Some preliminary investigations (relevant to the second order case of energy density  $H(\cdot, u, u', u'')$  when n = 1) had previously been performed via different methods by Aronsson and Aronsson-Barron in [A4, AB].

Apart from the intrinsic mathematical interest, the motivation to study higher order  $L^{\infty}$  minimisation problems comes from several diverse areas. In applied disciplines like Data Assimilation in the geosciences, PDE-constrained optimisation in aeronautics, etc. (see e.g. the model problem in [K9] and references therein, as well as the classical monograph [L]), a prevalent underlying problem is the construction of approximate solutions to second order ill-posed PDE problems. For instance, in the modelling of aquifers, one needs to solve a Poisson equation  $\Delta u = f$  coupled with a pointwise constraint of the form K(u) = k for given functions f, K, k. By minimising the error function  $|\Delta u - f|^2 + |K(u) - k|^2$  in  $L^{\infty}$ , one can obtain uniformly (absolutely) best approximations.

Minimisation problems in  $L^{\infty}$  similar to the above have also been studied in the context of differential geometry and in questions related to the Yamabe problem. In particular, the second author, together with Schwetlick [MS], and subsequently Sakellaris [S1] considered the problem of minimising the scalar curvature of a Riemannian metric on a given manifold and in a given conformal class. When formulated in terms of differential operators, this gives rise to a functional similar in structure to (1.1). This work uses different boundary conditions, however, and no attempt is made to prove uniqueness or find second order absolute minimisers. Nevertheless, some of the tools in the proofs of the above results originate in the above quoted papers.

We close with some remarks about generalised solutions to the equations governing the "extremals" of Calculus of Variations in  $L^{\infty}$ . In the scalar first order case, the theory of viscosity solutions of Crandall-Ishii-Lions (see [CIL, C, K8]) proved to be an apt framework within which the generally non-smooth solutions to the so-called Aronsson equation, which is a second order PDE, can be studied rigorously. However, viscosity solutions are of purely scalar nature and fail to work in either the vectorial or the higher order case (where we have either a second order system with discontinuous coefficients, or a fully nonlinear third order PDE). In the recent papers [K9, K10] a new theory of generalised solutions has been introduced which is based on a probabilistic representation of derivatives which do not exist classically and in the papers [AK, AyK, CKP, K11, K12, K13, KP, KP2, KP3] several results have been obtained in this framework. However, in the setting of the present paper, the particular structure of the problem at hand allows to prove directly existence of strong solutions to the fully nonlinear PDE (1.7).

#### 2. Proofs

In this section we establish the proofs of Theorems 2 and 3 and of Corollary 4.

**Proof of Theorem 2.** Fix  $\phi \in \mathcal{W}_0^{2,\infty}(\Omega)$  with  $\phi \not\equiv 0$  on  $\Omega$ . Since  $f_*$  is a harmonic function in  $L^1(\Omega)$ , it follows that

(2.1) 
$$\int_{\Omega} f_* \, \Delta \phi = 0.$$

We set

$$\Gamma_* := f_*^{-1}(\{0\})$$

By standard results on the nodal set of solutions to elliptic equations [HS] and the connectedness of  $\Omega$ , it follows that if  $f_* \neq 0$  then  $\Gamma_*$  is a Lebesgue nullset and if  $f_* \equiv 0$  then  $\Gamma_* = 0$ .

Let us first consider the case  $f_* \neq 0$ . Note that  $\Delta \phi$  cannot vanish almost everywhere on  $\Omega$  (as this would imply that  $\phi \equiv 0$  by uniqueness of solutions of the Dirichlet problem for the Laplace equation). Therefore, we deduce that  $f\Delta \phi \neq 0$ on a subset of positive Lebesgue measure in  $\Omega$ . Hence, (2.1) implies that there exist measurable sets  $\Omega^{\pm} \subseteq \Omega$  with  $\mathcal{L}^n(\Omega^{\pm}) > 0$  such that

$$\pm f_* \Delta \phi > 0$$
, a.e. on  $\Omega^{\pm}$ ,

where  $\mathcal{L}^n$  denotes the *n*-dimensional Lebesgue measure. If (1.6) holds true, then we have  $|\mathbf{F}(\cdot, \Delta u_*)| = e_*$  a.e. on  $\Omega$  and

$$\operatorname{sgn}(\Delta u_*) = \operatorname{sgn}(f_*) = \operatorname{sgn}(\Delta \phi),$$
 a.e. on  $\Omega^+$ .

As  $F(x, \cdot)$  is strictly increasing for a.e.  $x \in \Omega$ , this gives that

$$|\mathbf{F}(\cdot, \Delta u_* + \Delta \phi)| > |\mathbf{F}(\cdot, \Delta u_*)| = e_*, \text{ a.e. in } \Omega^+.$$

Therefore,

$$\mathcal{E}_{\infty}(u_*, \Omega) = e_* < \mathcal{E}_{\infty}(u_* + \phi, \Omega).$$

It remains to consider the case  $f_* \equiv 0$ . Then, the hypothesis (1.6) implies that  $\Delta u_* = 0$  almost everywhere and so  $E_{\infty}(u_*, \Omega) = 0$ . On the other hand, by arguing as above it follows that  $E_{\infty}(u_* + \phi, \Omega) > 0$  for any  $\phi \in \mathcal{W}_0^{2,\infty}(\Omega) \setminus \{0\}$ . Hence, we arrive at the same conclusion.

Finally, if  $\mathcal{O} \subseteq \Omega$  is a non-empty open set and  $\phi \in \mathcal{W}_0^{2,\infty}(\mathcal{O}) \setminus \{0\}$ , by repeating the previous reasoning with  $\mathcal{O}$  is the place of  $\Omega$  and  $\Gamma_* \cap \mathcal{O}$  in the place of  $\Gamma_*$ , we obtain once again the strict inequality  $\mathcal{E}_{\infty}(u_*, \mathcal{O}) < \mathcal{E}_{\infty}(u_* + \phi, \mathcal{O})$ . The theorem ensues.

The proof of Theorem 3 is more involved and requires some preparation. We begin with an elementary preliminary result.

**Lemma 5.** Let  $\Omega \subseteq \mathbb{R}^n$  be an open set and  $F \in C^2(\Omega \times \mathbb{R})$  a function satisfying (1.8). Then

(2.2) 
$$\operatorname{sgn}(\mathbf{F}(x,\xi)) = \operatorname{sgn}(\xi),$$

(2.3) 
$$\frac{1}{c}|\xi| \ge \left|\mathbf{F}(x,\xi)\right| \ge c|\xi|,$$

(2.4) 
$$(|\mathbf{F}(x,\cdot)|^p)_{\xi\xi}(\xi) \ge 0, \quad \text{if } p \ge \frac{1}{c^3} + 1,$$

(2.5) 
$$\mathbf{F}_{\xi}(x,\xi) |\xi| \geq c^2 |\mathbf{F}(x,\xi)|,$$

for all  $(x,\xi) \in \Omega \times \mathbb{R}$ , where c > 0 is the same constant as in (1.8).

**Proof of Lemma 5.** We first note that (2.2) is obvious, while (2.3) follows, by integration, from F(x,0) = 0 and the bounds  $c \leq F(x, \cdot) \leq 1/c$  of (1.8). For (2.4), we differentiate in  $\xi$  and use (1.8). Then

$$(|\mathbf{F}|^{p})_{\xi\xi} = p|\mathbf{F}|^{p-2} \Big( \mathbf{F} \mathbf{F}_{\xi\xi} + (p-1)(\mathbf{F}_{\xi})^{2} \Big)$$
  
 
$$\geq p|\mathbf{F}|^{p-2} \Big( -\frac{1}{c} + (p-1)c^{2} \Big),$$

which establishes the desired bound. Finally, (2.5) is a consequence of (1.8), which gives  $|\xi|F_{\xi}(x,\xi) \ge c|\xi|$ , and of (2.3), which gives  $c|\xi| \ge c^2|F(x,\xi)|$ .

We will construct the solutions to our problem by approximation with minimisers of  $L^p$  functionals. Therefore, we need to understand the behaviour of the latter.

**Proposition 6.** Suppose that  $F \in C^2(\Omega \times \mathbb{R})$  satisfies (1.8). Then for any  $p > c^{-3} + 1$  there exists a minimiser  $u_p$  of  $E_p$  over the space  $W^{2,p}_{u_0}(\Omega)$ . Moreover,  $u_p$  is a weak solution to the Dirichlet problem for the Euler-Lagrange equation associated with the functional  $E_p$ :

$$\begin{cases} \Delta \left( \left| \mathbf{F}(\cdot, \Delta u) \right|^{p-2} \mathbf{F}(\cdot, \Delta u) \mathbf{F}_{\xi}(\cdot, \Delta u) \right) = 0, & \text{in } \Omega, \\ u = u_0, & \text{on } \partial\Omega, \\ \mathbf{D}u = \mathbf{D}u_0, & \text{on } \partial\Omega. \end{cases} \end{cases}$$

Furthermore, there exist a (global) minimiser  $u_{\infty}$  of the functional  $E_{\infty}(\cdot, \Omega)$  over the space  $\mathcal{W}^{2,\infty}_{u_0}(\Omega)$  such that  $E_p(u_p) \longrightarrow E_{\infty}(u_{\infty}, \Omega)$  as  $p \to \infty$ . Also, there exists a subsequence  $(p_\ell)^{\infty}_1$  such that

$$\begin{cases} u_{p_{\ell}} \longrightarrow u_{\infty}, & \text{in } C^{1}(\overline{\Omega}), \\ D^{2}u_{p_{\ell}} \longrightarrow D^{2}u_{\infty}, & \text{in } L^{q}(\Omega, \mathbb{R}^{n \times n}), \text{ for all } q \in (1, \infty). \end{cases}$$

as  $\ell \to \infty$ .

**Proof of Proposition 6.** By (2.3)–(2.4) of Lemma 5, for  $p > c^{-3}+1$  the functional  $E_p$  is convex in  $W^{2,p}_{u_0}(\Omega)$  and

$$\mathbf{E}_p(u) \ge c(\mathcal{L}^n(\Omega))^{-1/p} \|\Delta u\|_{L^p(\Omega)},$$

for any  $u \in W^{2,p}_{u_0}(\Omega)$ . Since  $u - u_0 \in W^{2,p}_0(\Omega)$ , by the Calderon-Zygmund  $L^p$  estimates (e.g. [GT, GM]) and the Poincaré inequality we have a positive constant  $c_0 = c_0(p, \Omega)$  that

(2.6) 
$$\|\Delta u\|_{L^{p}(\Omega)} \geq c_{0} \|u\|_{W^{2,p}(\Omega)} - (c_{0}+1)\|u_{0}\|_{W^{2,p}(\Omega)}.$$

Hence,  $E_p$  is coercive on  $W^{2,p}_{u_0}(\Omega)$  and by setting

$$e_p := \inf \left\{ \mathcal{E}_p(u) : u \in W^{2,p}_{u_0}(\Omega) \right\}$$

we also have the bound

$$0 \le e_p \le \mathbf{E}_p(u_0) < \infty$$

because  $u_0 \in W^{2,p}(\Omega)$ . By applying the direct method of the Calculus of Variations (e.g. [D]), we deduce the existence of a global minimiser  $u_p \in W^{2,p}_{u_0}(\Omega)$ . Further,  $E_p$  is Gateaux differentiable at the minimiser as a result of the bound

$$\left| |\mathbf{F}(\cdot,\xi)|^{p-1} \mathbf{F}_{\xi}(\cdot,\xi) \right| \leq C |\xi|^{p-1}$$

and well-known results (see e.g. [D, GM]).

Consider a family of minimisers  $(u_p)_{p \ge p_0}$  where

$$p_0 := \{ \text{integer part of } \max\{n, c^{-3}\} + 1 \}$$

and fix  $k \in \mathbb{N}$ . For any  $p \ge k$ , by (2.3), Hölder's inequality and the minimality we have

(2.7) 
$$c \|\Delta u_p\|_{L^k(\Omega)} (\mathcal{L}^n(\Omega))^{-1/k} \leq \mathbf{E}_k(u_p) \leq \mathbf{E}_p(u_p) \leq \mathbf{E}_p(u_0) \leq \mathbf{E}_\infty(u_0,\Omega)$$

and hence  $(\Delta u_p)_{p \ge p_0}$  is bounded in  $L^k(\Omega)$ . By the previous arguments and (2.6), we conclude that  $(u_p)_{p \ge p_0}$  is bounded in  $W^{2,k}_{u_0}(\Omega)$  for any  $k \in \mathbb{N}$ . By a standard diagonal argument, weak compactness and the Morrey theorem, there exists

$$u_{\infty} \in \bigcap_{1 < k < \infty} W^{2,\infty}_{u_0}(\Omega)$$

such that the desired convergences hold true along a subsequence as  $p_{\ell} \to \infty$ . When we pass to the limit as  $\ell \to \infty$  in (2.7), the weak lower semicontinuity of the  $L^k$ norm implies

$$\|\Delta u_{\infty}\|_{L^{k}(\Omega)} \leq \frac{(\mathcal{L}^{n}(\Omega))^{1/k}}{c} \operatorname{E}_{\infty}(u_{0}, \Omega).$$

Letting  $k \to \infty$  we obtain  $\Delta u_{\infty} \in L^{\infty}(\Omega)$ . Thus,  $u_{\infty} \in \mathcal{W}^{2,\infty}_{u_0}(\Omega)$ , as desired. It remains to show the convergence of  $\mathcal{E}_p(u_p)$  and minimality of  $u_{\infty}$ .

Hölder's inequality and minimality show that

$$E_p(u_p) \le E_p(u_q) \le E_q(u_q)$$
, whenever  $p \le q$ .

Therefore, the limit  $\lim_{p\to\infty} E_p(u_p)$  exists. Since  $u_p - u_\infty \in W_0^{2,p}(\Omega)$ , for any  $\phi \in W_0^{2,\infty}(\Omega)$  the minimality and Hölder's inequality imply

(2.8)  

$$E_{\infty}(u_{\infty}, \Omega) = \lim_{k \to \infty} E_{k}(u_{\infty})$$

$$\leq \liminf_{k \to \infty} \left(\liminf_{\ell \to \infty} E_{k}(u_{p_{\ell}})\right)$$

$$\leq \lim_{p \to \infty} E_{p}(u_{p})$$

$$\leq \limsup_{p \to \infty} E_{p}(u_{\infty} + \phi)$$

$$\leq E_{\infty}(u_{\infty} + \phi, \Omega).$$

Inequality (2.8) implies that  $u_{\infty}$  is indeed a global minimiser of  $E_{\infty}(\cdot, \Omega)$  over  $\mathcal{W}^{2,\infty}_{u_{0}}(\Omega)$ . In addition, the choice  $\phi = 0$  in (2.8) gives

$$\mathbf{E}_{\infty}(u_{\infty},\Omega) \leq \lim_{p \to \infty} \mathbf{E}_p(u_p) \leq \mathbf{E}_{\infty}(u_{\infty},\Omega)$$

and hence  $\mathbf{E}_p(u_p) \longrightarrow \mathbf{E}_{\infty}(u_{\infty}, \Omega)$ , as claimed.

The next result is an essential part of our constructions and this is the only point at which we make use of the  $C^{1,1}$  boundary regularity of  $\partial\Omega$  and the slightly higher regularity of the boundary condition  $u_0 \in \mathcal{W}^{2,\infty}(\Omega)$  (that is, that the Laplacian  $\Delta u_0$  is continuous on  $\overline{\Omega}$ ).

Subsequently, we will be using the following symbolisation for the *r*-neighbourhood of the boundary  $\partial \Omega$  in  $\Omega$ :

$$\Omega_r := \{ x \in \Omega : \operatorname{dist}(x, \partial \Omega) < r \},\$$

for r > 0.

**Lemma 7** (Improving the boundary data). Let  $\Omega \subseteq \mathbb{R}^n$  be a bounded open set with  $C^{1,1}$  boundary and consider a function  $u_0 \in \mathcal{W}^{2,\infty}(\Omega)$  with  $\Delta u_0 \in C(\overline{\Omega})$ . Then, for any  $\epsilon > 0$  there exist a number  $r = r(\varepsilon) > 0$  and a function  $w = w(\varepsilon) \in \mathcal{W}^{2,\infty}_{u_0}(\Omega)$  such that

$$\|\Delta w\|_{L^{\infty}(\Omega_r)} \leq \varepsilon.$$

In other words, given any boundary condition  $u_0 \in \mathcal{W}^{2,\infty}(\Omega)$  such that  $\Delta u_0$  is continuous up to the boundary, we can find another function w in the same space with the same boundary data, the Laplacian of which is as small as desired in a neighbourhood of the boundary.

**Proof of Lemma 7.** Since  $\partial\Omega$  is  $C^{1,1}$ -regular, the result of the appendix establishes that the distance function  $\operatorname{dist}(\cdot, \partial\Omega)$  belongs to  $C^{1,1}(\overline{\Omega}_{2r_0})$  for some  $r_0 > 0$  small enough.

Let d be an extension of dist $(\cdot, \partial \Omega)$  from  $\Omega_{r_0}$  to  $\overline{\Omega}$  which is in the space  $W^{2,\infty}(\Omega)$ . Extend  $\Delta u_0$  and d by zero on  $\mathbb{R}^n \setminus \Omega$ . Let  $(\eta^{\delta})_{\delta>0} \subseteq C_c^{\infty}(\mathbb{R}^n)$  be a standard mollifying family (as e.g. in [E]). We set

(2.9) 
$$v_{\delta} := u_0 - \frac{d^2}{2} \left( \eta^{\delta} * \Delta u_0 \right)$$

Then  $v_{\delta} - u_0 \in \mathcal{W}^{2,\infty}_0(\Omega)$  since d = 0 on  $\partial \Omega$  and

$$Dv_{\delta} = Du_{0} - d\left\{\frac{d}{2}D(\eta^{\delta} * \Delta u_{0}) + (\eta^{\delta} * \Delta u_{0})Dd\right\},$$
  

$$D^{2}v_{\delta} = D^{2}u_{0} - (Dd \otimes Dd)(\eta^{\delta} * \Delta u_{0}) - d\left\{\frac{d}{2}D^{2}(\eta^{\delta} * \Delta u_{0}) + D(\eta^{\delta} * \Delta u_{0}) \otimes Dd + Dd \otimes D(\eta^{\delta} * \Delta u_{0}) + (\eta^{\delta} * \Delta u_{0})D^{2}d\right\}$$

By using that

$$\operatorname{tr}(\mathrm{D}d\otimes\mathrm{D}d) = |\mathrm{D}d|^2 = 1 \text{ on } \Omega_{r_0}$$

for  $0 < r < r_0$  we deduce

(2.10)  
$$\begin{aligned} \|\Delta v_{\delta}\|_{L^{\infty}(\Omega_{r})} &\leq \left\|\Delta u_{0} - \eta^{\delta} * \Delta u_{0}\right\|_{L^{\infty}(\Omega_{r_{0}})} + C_{1}r\left(\left\|\Delta u_{0} * \mathrm{D}^{2}\eta^{\delta}\right\|_{L^{\infty}(\mathbb{R}^{n})} \\ &+ \left\|\Delta u_{0} * \mathrm{D}\eta^{\delta}\right\|_{L^{\infty}(\mathbb{R}^{n})} + \left\|\Delta u_{0}\right\|_{L^{\infty}(\mathbb{R}^{n})}\right)\end{aligned}$$

for some constant  $C_1 > 0$ . Since

$$\eta^{\delta}(x) = \delta^{-n} \eta(|x|/\delta)$$

for some fixed function  $\eta \in C_c^{\infty}(\mathbb{B}_1(0))$ , by Young's inequality for convolutions we have the estimate

(2.11) 
$$\left\|\Delta u_0 * \mathbf{D}^k \eta^\delta\right\|_{L^\infty(\mathbb{R}^n)} \le \frac{1}{\delta^k} \left\|\Delta u_0\right\|_{L^\infty(\mathbb{R}^n)} \left\|\mathbf{D}^k \eta\right\|_{L^1(\mathbb{R}^n)}, \quad k = 1, 2.$$

Hence, by invoking (2.11) we see that (2.10) gives

$$\|\Delta v_{\delta}\|_{L^{\infty}(\Omega_r)} \leq \|\Delta u_0 - \eta^{\delta} * \Delta u_0\|_{L^{\infty}(\Omega_{r_0})} + C_2 r \left(\frac{1}{\delta^2} + \frac{1}{\delta} + 1\right) \|\Delta u_0\|_{L^{\infty}(\Omega_{r_0})},$$

where  $C_2 = C_1 \|\eta\|_{W^{2,1}(\mathbb{R}^n)}$ . By choosing

$$\delta := r^{1/4}$$

and also choosing  $r_0$  sufficiently small, we obtain the desired statement as a consequence of the continuity of  $\Delta u_0$ .

Now we can show that the minimiser  $u_{\infty}$  obtained in Proposition 6 satisfies the desired formula of part (I) in Theorem 3. The rest of the proof is then not difficult.

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**Proof of Theorem 3**. Let us begin by setting

$$e_{\infty} := \inf \left\{ \mathcal{E}_{\infty}(u, \Omega) : u \in \mathcal{W}^{2, \infty}_{u_0}(\Omega) \right\}.$$

If  $e_{\infty} = 0$ , then everything in Theorem 3 follows from Proposition 6 or is trivial. (Note that in this case  $e_p = 0$  for every p and every minimiser of the corresponding functionals is a harmonic function.) Therefore, we may assume that  $e_{\infty} > 0$ .

Since  $e_p \longrightarrow e_{\infty}$  as  $p \to \infty$  by Proposition 6, it follows that  $e_p > 0$  for large p, say for  $p \ge p_0$ . Then for all  $p \ge p_0$ , the formula in (IV) gives rise to a measurable function  $f_p : \Omega \longrightarrow \mathbb{R}$ . Then the Euler-Lagrange equation in Proposition 6 can be expressed in the form

(2.12) 
$$\Delta f_p = 0, \quad \text{on } \Omega.$$

That is,  $f_p$  is harmonic on  $\Omega$  and hence belongs to  $C^{\infty}(\Omega)$ .

Let p' = p/(p-1) be the conjugate exponent of  $p \in (1, \infty)$ . Then (1.8) implies

$$\left( \oint_{\Omega} |f_p|^{p'} \right)^{1/p'} = \frac{1}{e_p^{p-1}} \left( \oint_{\Omega} \left| \mathbf{F}^{p-1}(\cdot, \Delta u_p) \mathbf{F}_{\xi}(\cdot, \Delta u_p) \right|^{p'} \right)^{1/p'}$$
$$\leq \frac{1}{c e_p^{p-1}} \left( \oint_{\Omega} \left| \mathbf{F}^{p-1}(\cdot, \Delta u_p) \right|^{p/(p-1)} \right)^{(p-1)/p}$$
$$= \frac{1}{c},$$

which gives the following uniform  $L^1$  bound of  $(f_p)_{p \ge p_0} \subseteq C^{\infty}(\Omega)$ :

(2.13) 
$$\|f_p\|_{L^1(\Omega)} \leq \mathcal{L}^n(\Omega) \left( \oint_{\Omega} |f_p|^{p'} \right)^{1/p'} \leq \frac{\mathcal{L}^n(\Omega)}{c}$$

By the mean value theorem for harmonic functions and by the standard interior derivative estimates (e.g. [GT]) we have, for any  $k \in \mathbb{N} \cup \{0\}$  and for any compactly contained  $\mathcal{O} \subseteq \Omega$ , a constant  $C = C(k, \mathcal{O}, \Omega) > 0$  such that

$$\|\mathbf{D}^k f_p\|_{L^{\infty}(\mathcal{O})} \leq C \|f_p\|_{L^1(\Omega)}.$$

Hence, the family  $(f_p)_{p\geq p_0}$  is bounded (in the locally convex sense) in the topology of  $C^{\infty}(\Omega)$  and as a consequence there exist  $f_{\infty} \in C^{\infty}(\Omega)$  and a sequence  $p_{\ell} \to \infty$ such that  $f_{p_{\ell}} \longrightarrow f_{\infty}$  as  $\ell \to \infty$ . From (2.12) it follows that  $f_{\infty}$  is harmonic:

(2.14) 
$$\Delta f_{\infty} = 0, \quad \text{on } \Omega.$$

Fix r > 0 and consider the inner *r*-neighbourhood  $\Omega_r$  of  $\partial\Omega$ . Since  $f_{p_\ell} \longrightarrow f_\infty$  in  $C(\overline{\Omega \setminus \Omega_r})$ , inequality (2.13) implies that

$$\|f_{\infty}\|_{L^{1}(\Omega\setminus\Omega_{r})} = \lim_{\ell\to\infty} \|f_{p_{\ell}}\|_{L^{1}(\Omega\setminus\Omega_{r})} \leq \frac{\mathcal{L}^{n}(\Omega)}{c}.$$

Letting  $r \to 0$  we conclude that  $f_{\infty} \in L^1(\Omega)$  and

(2.15) 
$$\|f_{\infty}\|_{L^{1}(\Omega)} \leq \frac{\mathcal{L}^{n}(\Omega)}{c}$$

We now show that  $f_{\infty} \neq 0$  using Lemma 7. Fix  $\epsilon > 0$  small and let r > 0 and  $w \in \mathcal{W}^{2,\infty}_{u_0}(\Omega)$  be as constructed in Lemma 7 with  $|\Delta w| \leq \varepsilon$  on the inner zone

 $\Omega_r \subseteq \Omega$  of the boundary  $\partial \Omega$ . Since  $u_p - w \in W_0^{2,p}(\Omega)$ , it is an admissible test function and by (2.12), integration by parts gives

$$\int_{\Omega} f_p \,\Delta(u_p - w) \,=\, 0.$$

Hence, by the above together with (1.9), (1.8), (2.2) and (2.5), we obtain

$$\begin{split} \int_{\Omega} f_p \,\Delta w &= \int_{\Omega} f_p \,\Delta u_p \\ &= \frac{1}{e_p^{p-1}} \int_{\Omega} \left| \mathbf{F}(\cdot, \Delta u_p) \right|^{p-2} \mathbf{F}(\cdot, \Delta u_p) \,\mathbf{F}_{\xi}(\cdot, \Delta u_p) \Delta u_p \\ &= \frac{1}{e_p^{p-1}} \int_{\Omega} \left| \mathbf{F}(\cdot, \Delta u_p) \right|^{p-1} \mathbf{F}_{\xi}(\cdot, \Delta u_p) \,|\Delta u_p| \\ &\geq \frac{c^2}{e_p^{p-1}} \int_{\Omega} \left| \mathbf{F}(\cdot, \Delta u_p) \right|^p. \end{split}$$

Thus,

(2.16) 
$$\int_{\Omega} f_p \,\Delta w \geq c^2 \,\mathcal{L}^n(\Omega) \,e_p.$$

Now, we use (2.16), (2.13) and Lemma 7 to estimate

$$c^{2} \mathcal{L}^{n}(\Omega) e_{p} \leq \int_{\Omega_{r}} f_{p} \Delta w + \int_{\Omega \setminus \Omega_{r}} f_{p} \Delta w$$
$$\leq \varepsilon \|f_{p}\|_{L^{1}(\Omega)} + \int_{\Omega \setminus \Omega_{r}} f_{p} \Delta w$$
$$\leq \frac{\varepsilon \mathcal{L}^{n}(\Omega)}{c} + \int_{\Omega \setminus \Omega_{r}} f_{p} \Delta w.$$

Recalling that  $e_{p_{\ell}} \longrightarrow e_{\infty}$  and also  $f_{p_{\ell}} \longrightarrow f_{\infty}$  in  $C(\overline{\Omega \setminus \Omega_r})$ , we pass to the limit as  $\ell \to \infty$  to find that

$$\int_{\Omega \setminus \Omega_r} f_{\infty} \Delta w \ge \mathcal{L}^n(\Omega) \Big( c^2 e_{\infty} - \frac{\varepsilon}{c} \Big).$$

By choosing  $\varepsilon > 0$  small enough, we deduce that

$$\int_{\Omega \setminus \Omega_r} f_\infty \, \Delta w \, > \, 0,$$

which implies  $f_{\infty} \neq 0$ , as claimed.

Let us now define the map

$$\Phi \ : \ \Omega \times \mathbb{R} \longrightarrow \Omega \times \mathbb{R}, \quad \Phi(x,\xi) \ := \ (x, \mathcal{F}(x,\xi)).$$

Under the assumption (1.8), this is a  $C^2$  diffeomorphism.

Since the inverse function of  $t \mapsto |t|^{p-2}t$  is given by  $s \mapsto \operatorname{sgn}(s)|s|^{1/(p-1)}$ , we may rewrite the formula (1.9) defining the harmonic function  $f_p$ , as

$$\mathbf{F}(\cdot, \Delta u_p) = e_p |f_p|^{\frac{1}{p-1}} \left[ \mathbf{F}_{\xi}(\cdot, \Delta u_p) \right]^{-\frac{1}{p-1}} \operatorname{sgn}(f_p)$$

or as

(2.17) 
$$(x, \Delta u_p(x)) = \Phi^{-1} \left( \cdot, e_p |f_p|^{\frac{1}{p-1}} \left[ F_{\xi}(\cdot, \Delta u_p) \right]^{-\frac{1}{p-1}} \operatorname{sgn}(f_p) \right) (x),$$

for  $x \in \Omega \setminus \Gamma_{\infty}$ . On any compact set  $K \subseteq \Omega \setminus \Gamma_{\infty}$ , we have the uniform convergence  $f_{p_{\ell}} \longrightarrow f_{\infty}$  as  $\ell \to \infty$ , whereas  $F_{\xi}$  is uniformly bounded from above and below by (1.8). Hence, by restricting ourselves along the subsequence  $p_{\ell}$  and letting  $\ell \to \infty$  we obtain uniform convergence of the right-hand side of (2.17) to  $\Phi^{-1}(\cdot, e_{\infty} \operatorname{sgn}(f_{\infty}))$  on K. But since we already know that  $\Delta u_p \longrightarrow \Delta u_{\infty}$  weakly in  $L^2(\Omega)$ , it follows that

(2.18) 
$$(x, \Delta u_{\infty}(x)) = \Phi^{-1} \Big( x, e_{\infty} \operatorname{sgn} \big( f_{\infty}(x) \big) \Big), \quad x \in K.$$

As a consequence,

$$F(x, \Delta u_{\infty}(x)) = e_{\infty} \operatorname{sgn}(f_{\infty}(x)), \quad x \in K.$$

Now let us recall that  $\mathcal{L}^n(\Gamma_{\infty}) = 0$ . This is a consequence of general regularity results for nodal sets of solution to elliptic equations [HS]. The statement of item (I) then follows.

In order to prove item (II), we note that (2.18) implies that  $\Delta u_{\infty} \in C^2(\Omega \setminus \Gamma_{\infty})$ . The desired statement then follows from standard Schauder theory [GT].

For item (III), first recall the subsequential convergence of the  $E_p$ -minimisers  $(u_p)_1^{\infty}$  of Proposition 6 along  $(p_{\ell})_1^{\infty}$  as  $\ell \to \infty$ . We also have the desired respective convergence of the global infima  $(e_p)_1^{\infty}$  of the energies.

The a.e. convergence of the Laplacians  $(\Delta u_{p_{\ell}})_{\ell=1}^{\infty}$  follows from the fact that  $\Omega \setminus \Gamma_{\infty}$  has full Lebesgue measure and that the sequence converges locally uniformly thereon.

The strong convergence of the Laplacians  $(\Delta u_{p_{\ell}})_{\ell=1}^{\infty}$  in  $L^{q}(\Omega)$  for all  $q \in (1, \infty)$  is a consequence of the Vitaly convergence theorem (see e.g. [FL]) and of the following facts:

i) the weak convergence of the Laplacians over the same spaces,

ii) the a.e. convergence of the Laplacians on  $\Omega$ ,

- iii) the boundedness of  $\Omega$ ,
- iv) the  $L^q$  equi-integrability estimate

$$\|\Delta u_{p_{\ell}}\|_{L^{q}(E)} \leq \left(\sup_{\ell \in \mathbb{N}} \|\Delta u_{p_{\ell}}\|_{L^{q+1}(\Omega)}\right) \left(\mathcal{L}^{n}(E)\right)^{\frac{1}{q(q+1)}}$$

which holds true for any measurable subset  $E \subseteq \Omega$ .

Finally, the statement in part (IV) has already been proven.

We conclude this section by noting that Corollary 4 is an immediate consequence of Theorem 2, Theorem 3, and the observation that  $\mathcal{L}^n(\Gamma_{\infty}) = 0$  if  $e_{\infty} > 0$ . On the other hand,  $u_{\infty}$  is a harmonic function if  $e_{\infty} = 0$  and the result follows trivially.

Appendix:  $C^{1,1}$  regularity of the distance function for  $C^{1,1}$  domains

Suppose that  $\Omega \subseteq \mathbb{R}^n$  is a bounded open set with  $C^{1,1}$  boundary  $\partial \Omega$ . Let

$$d \equiv \operatorname{dist}(\cdot, \partial \Omega) : \Omega \longrightarrow \mathbb{R}$$

symbolise the distance function to the boundary. For r > 0, let  $\Omega_r$  denote again the inner *r*-neighbourhood of  $\partial \Omega$  in  $\Omega$ :

$$\Omega_r = \{ x \in \Omega : d(x) < r \}.$$

In this appendix we establish that  $d \in C^{1,1}(\overline{\Omega_r})$  when r is sufficiently small. This fact is probably known but we couldn't locate a precise reference in the literature.

The  $C^2$  regularity of the distance function for a  $C^2$  boundary  $\partial\Omega$  is a classical result, see e.g. [GT, Appendix 14.6]. On the other hand, the case of  $C^1$  regularity of the distance function when the boundary  $\partial\Omega$  is  $C^1$  holds under the extra hypothesis that the distance is realised at *one* point; see e.g. [F].

In order to prove the desired  $C^{1,1}$  regularity of the distance function near  $\partial\Omega$ when the boundary itself is a  $C^{1,1}$  manifold (which we utilised in Lemma 7), we first note the following fact: suppose that r > 0 is such that 1/r is larger than the essential supremum of the curvature of  $\partial\Omega$ . If  $x \in \Omega$  and  $y \in \partial\Omega$  with  $|x-y| = s \leq r$ and such that the tangent hyperplanes of  $\partial\Omega$  and  $\partial\mathbb{B}_s(x)$  coincide at y, then it follows that  $\mathbb{B}_s(x) \subseteq \Omega$  and  $\partial\mathbb{B}_s(x) \cap \partial\Omega = \{y\}$ . (This is easy to see when  $\partial\Omega$  is  $C^2$  regular and follows by an approximation of  $\partial\Omega$  with  $C^2$  manifolds otherwise.) Therefore, in the above situation, it follows that d(x) = s. Moreover, for  $x \in \Omega$  with  $d(x) \leq r$ , it follows that there exists a *unique* point  $y \in \partial\Omega$  such that |x-y| = d(x). Moreover, if  $\nu$  denotes the outer normal vector on  $\partial\Omega$ , then  $x = y - d(x)\nu(y)$  and  $Dd(x) = -\nu(y)$ .

Now fix  $x_0 \in \partial \Omega$ . Our aim is to prove  $C^{1,1}$  regularity of d near  $x_0$ . To this end, we may assume without loss of generality that there exist open sets  $U \subseteq \mathbb{R}^n$  and  $V \subseteq \mathbb{R}^{n-1}$  such that  $x_0 \in U$  and  $0 \in V$  and there exists a function  $f \in C^{1,1}(V)$ such that

$$\Omega \cap U = \left\{ (x', x_n) \in V \times \mathbb{R} : x_n > f(x') \right\} \cap U$$

and such that  $x_0 = (0, f(0))$  and Df(0) = 0. Define

$$N(x') := \frac{(-Df(x'), 1)}{\sqrt{1 + |Df(x')|^2}}, \quad x' \in V,$$

so that  $N(x') = -\nu(x', f(x'))$  for  $x' \in V$ . Note that this map is Lipschitz continuous. We now define  $\Psi : V \times \mathbb{R} \longrightarrow \mathbb{R}^n$  by

$$\Psi(x',t) := (x', f(x')) + t N(x').$$

Then  $\Psi$  is  $C^{0,1}$  near (0,0). Note also that  $\Psi$  is injective in a sufficiently small neighbourhood of (0,0): if we had  $\Psi(x',s) = \Psi(y',t) =: z$  with  $0 < s \leq t \leq r$ , then it would follow that  $\partial \mathbb{B}_t(z)$  and  $\partial \Omega$  have the same tangent hyperplanes at (y', f(y')). By the above observations, this would imply that  $\mathbb{B}_t(z) \subseteq \Omega$  and  $\partial \mathbb{B}_t(z) \cap \partial \Omega = \{(y', f(y'))\}$ , and therefore x' = y' and s = t. So if U and V are chosen appropriately, then  $\Psi$  is a bijection between  $V \times [0, r)$  and  $U \cap \overline{\Omega}$ . Moreover, we compute

$$D_j \Psi_i(x',t) = \delta_{ij} + t D_j N_i(x'), \quad i,j = 1,..., n-1,$$

and

$$D_j \Psi_n(x',t) = D_j f(x') + t D_j N_n(x'), \qquad j = 1, \dots, n-1,$$

while

$$D_t \Psi(x',t) = N(x').$$

Since Df(0) = 0 and N(0) = (0, ..., 0, 1), it follows that  $D\Psi$  is of full rank in some neighbourhood of (0, 0) and moreover, the inverse  $(D\Psi)^{-1}$  is essentially bounded in this neighbourhood. That is, by making V and r smaller if necessary, without loss of generality we may assume that

$$\Psi^{-1} \in C^{0,1}(U \cap \overline{\Omega}; V \times [0, r)).$$

Now note that

$$Dd(\Psi(x',t)) = N(x')$$

whenever t > 0 is small enough. Hence if  $\pi$  denotes the projection onto  $\mathbb{R}^{n-1} \times \{0\}$ , then we obtain the formula

$$\mathrm{D}d = N \circ \pi \circ \Psi^{-1}$$

near  $x_0$ . The right-hand side is of class  $C^{0,1}$ , and thus d is of class  $C^{1,1}$  near  $x_0$ . A compactness argument then proves the above statement.

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